

Chapter 5

Applications of Statistical Thermodynamics

5.1 Entropy S

M. C. E.

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Macroscopic parameters: $E, \{x_i\}$

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$$S = S(E, \{x_i\}) \quad (5.1)$$

$$= k_B \ln \Omega(E, \{x_i\}) \quad (5.2)$$

$$dS = \left(\frac{\partial S}{\partial E} \right)_{\{x_i\}} dE + \sum_i \left(\frac{\partial S}{\partial x_i} \right)_{E, \{x_i\}'} \cdot dx_i \quad (5.3)$$

$$= \frac{1}{T} dE + \sum_i \frac{\bar{X}_i}{T} dx_i \quad (5.4)$$

∴ one can obtain the other macroscopic quantities T & $\{\bar{X}_i\}$ from its derivatives.

e. g. 1 Gas system

$$x = V, \bar{X} = P \quad (5.5)$$

$$S = S(E, V), \frac{\partial S}{\partial E} = \frac{1}{T}, \frac{\partial S}{\partial V} = \frac{P}{T} \quad (5.6)$$

e. g. 2 Magnetic system

$$x = B \text{ (external magnetic field)} \quad (5.7)$$

$$\bar{X} = M \text{ (magnetization)} \quad (5.8)$$

$$S = S(E, B) \quad (5.9)$$

$$\frac{\partial S}{\partial E} = \frac{1}{T}, \quad \frac{\partial S}{\partial B} = \frac{M}{T} \quad (5.10)$$

e. g. 1 & e. g. 2 \longrightarrow \exists four macroscopic quantities

$$\{(E, T), (V, P)\}, \{(E, T), (B, M)\} \quad (5.11)$$

\downarrow

Only two of the four quantities are indep.

$\therefore \exists$ two conditions

$$\left(\frac{\partial S}{\partial E}\right)_x = \frac{1}{T}, \quad \left(\frac{\partial S}{\partial x}\right)_E = \frac{\bar{X}}{T} \quad (5.12)$$

\downarrow

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$$E = E(T, x) \quad \bar{X} = \bar{X}(T, x) \quad (5.13)$$

5.2 Thermodynamic Potentials and Maxwell's relations

A thermodynamic function whose dimension is energy is called a “thermodynamic potential.”

5.2.1 Internal Energy U

$$U = U(S, x) \quad (5.14)$$

$$dU = \left(\frac{\partial U}{\partial S}\right)_x dS + \left(\frac{\partial U}{\partial x}\right)_S \cdot dx \quad (5.15)$$

$$dS = \frac{1}{T}dU + \frac{\bar{X}}{T}dx \longrightarrow dU = T dS - \bar{X} \cdot dx \quad (5.16)$$

$$\therefore \left(\frac{\partial U}{\partial S}\right)_x = T, \quad \left(\frac{\partial U}{\partial x}\right)_S = -\bar{X} \quad (5.17)$$

$$\{(S, T) (x, \bar{X})\} \quad (5.18)$$

$$\frac{\partial^2 U}{\partial x \partial S} = \frac{\partial^2 U}{\partial S \partial x} \longrightarrow \left(\frac{\partial T}{\partial x} \right)_S = - \left(\frac{\partial \bar{X}}{\partial S} \right)_x \quad (5.19)$$

Maxwell's relation

∃ four thermodynamic variables

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Only two of the four variables: indep.

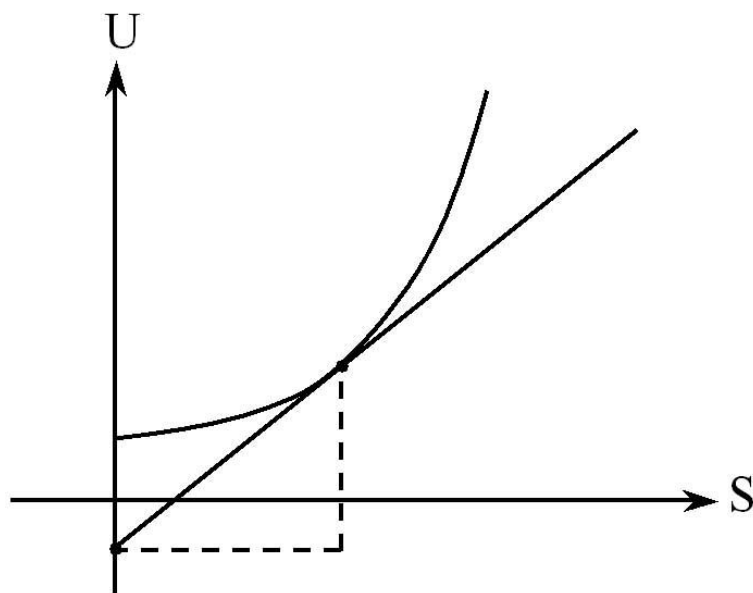
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The other two quantities can be obtained from the partial derivatives of the thermodynamic potential with respect to the chosen indep. thermodynamic variables

5.2.2 Helmholtz Free Energy

$$F = F(T, x) \longleftarrow U = U(S, x) \quad (5.20)$$

Legendry transformation



$$\text{Slope of the tangent: } T \rightarrow \frac{\partial U}{\partial S} = T \quad (5.21)$$

$$\text{Intercept of the tangent: } F \quad (5.22)$$

$$\longrightarrow T = \frac{U - F}{S} \quad (5.23)$$

$$\begin{aligned} \therefore F(T, x) = U(S, x) - TS, \quad \left(\frac{\partial U}{\partial S}\right)_x &= T \\ S &= S(T, x) \end{aligned} \quad (5.24)$$

$$dF = dU - T dS - S dT \quad (5.25)$$

$$= -S dT - \bar{X} \cdot dx \quad (5.26)$$

$$\therefore \left(\frac{\partial F}{\partial T}\right)_x = -S, \quad \left(\frac{\partial F}{\partial x}\right)_T = -\bar{X}. \quad (5.27)$$

$$\text{Maxwell's relation: } \left(\frac{\partial S}{\partial x}\right)_T = \left(\frac{\partial \bar{X}}{\partial T}\right)_x$$

5.2.3 Enthalpy H

$$U = U(S, x) \longrightarrow H = H(S, \bar{X}) \quad (5.28)$$

$$\left(\frac{\partial U}{\partial x}\right)_S = \frac{U - H}{x} = -\bar{X} \quad (5.29)$$

$$\therefore H = U - \left(\frac{\partial U}{\partial x}\right)_S \cdot x \quad (5.30)$$

$$= U + \bar{X} \cdot x, \quad \bar{X} = -\left(\frac{\partial U}{\partial x}\right)_S \quad (5.31)$$

↓

$$x = x(S, \bar{X})$$

$$dH = dU + \bar{X} \cdot dx + x \cdot d\bar{X} \quad (5.32)$$

$$= T \cdot dS + x \cdot d\bar{X} \quad (5.33)$$

$$\therefore \left(\frac{\partial H}{\partial S} \right) = T, \quad \left(\frac{\partial H}{\partial \bar{X}} \right) = x. \quad (5.34)$$

$$\text{Maxwell's relation: } \left(\frac{\partial T}{\partial \bar{X}} \right)_S = \left(\frac{\partial x}{\partial S} \right)_{\bar{X}}$$

5.2.4 Gibbs Free Energy G

$$F = F(T, x) \longrightarrow G = G(T, \bar{X}) \quad (5.35)$$

$$\left(\frac{\partial F}{\partial x} \right)_T = \frac{F - G}{x} = -\bar{X} \quad (5.36)$$

$$\therefore G = F - \left(\frac{\partial F}{\partial x} \right)_T \cdot x = F + \bar{X} \cdot x \quad (5.37)$$

$$dG = dF + \bar{X} \cdot dx + x \cdot d\bar{X} \quad (5.38)$$

$$= -S dT + x \cdot d\bar{X} \quad (5.39)$$

$$\therefore \left(\frac{\partial G}{\partial T} \right)_x = -S, \quad \left(\frac{\partial G}{\partial \bar{X}} \right)_T = x. \quad (5.40)$$

$$\text{Maxwell's relation: } - \left(\frac{\partial S}{\partial \bar{X}} \right)_T = \left(\frac{\partial x}{\partial T} \right)_{\bar{X}}$$

5.3 Extensive and Intensive Parameters

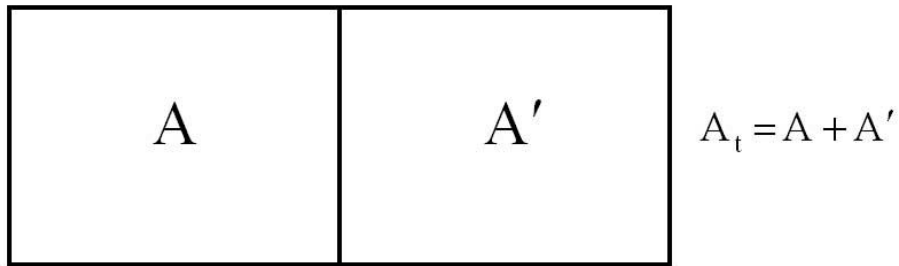
Let us consider a system A_t whose macrostate is specified by macroscopic parameters $(E, \{x_i\}, N)$.

↓

S

↓

$T, \{\bar{X}_i\}$



$$N_t = N + N', \quad E_t = E + E', \quad S_t = S + S', \quad (5.41)$$

$$V_t = V + V', \quad M_t = M + M', \quad (5.42)$$

$$T_t = T = T', \quad P_t = P = P', \quad B_t = B + B' \quad (5.43)$$

y : a macroscopic parameter

$$y_t = y + y' \implies y : \text{extensive parameter}$$

$$y_t = y = y' \implies y : \text{intensive parameter}$$

extensive parameter \longrightarrow intensive parameter

$$\text{int.} = \frac{\text{ext.}}{\text{ext.}} \quad (5.44)$$

$$\text{e. g. } e = \frac{E}{N}, \quad m = \frac{M}{N}, \quad \dots \quad (5.45)$$

5.4 Relations between partial derivatives of several variables

There are three variables (x, y, z) .

But, only two variables are independent.

$$\text{e. g. } (x, y) : \text{indep.}, \quad z = z(x, y) \quad (5.46)$$

$$dz = \left(\frac{\partial z}{\partial x} \right)_y dx + \left(\frac{\partial z}{\partial y} \right)_x dy \quad (5.47)$$

$$\left(\frac{\partial y}{\partial x}\right)_z \quad ? \quad (5.48)$$

$$\left(\frac{\partial y}{\partial x}\right)_z = -\frac{\left(\frac{\partial z}{\partial x}\right)_y}{\left(\frac{\partial z}{\partial y}\right)_x} : \text{Euler's relation} \quad (5.49)$$

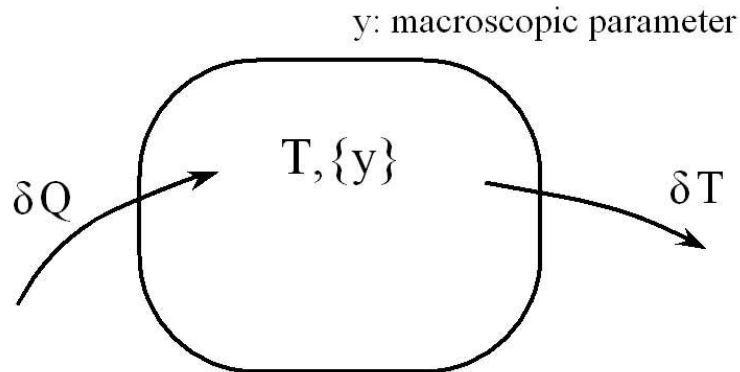
$$F = F(x, y(x, z)), \quad y = y(x, z) \quad (5.50)$$

$$dF = \left(\frac{\partial F}{\partial x}\right)_y dx + \left(\frac{\partial F}{\partial y}\right)_x dy \quad (5.51)$$

$$\left(\frac{\partial F}{\partial x}\right)_z = \left(\frac{\partial F}{\partial x}\right)_y + \left(\frac{\partial F}{\partial y}\right)_x \cdot \left(\frac{\partial y}{\partial x}\right)_z \quad (5.52)$$

$$\left(\frac{\partial F}{\partial z}\right)_x = \left(\frac{\partial F}{\partial y}\right)_x \cdot \left(\frac{\partial y}{\partial z}\right)_x \quad (5.53)$$

5.5 Specific Heat and Susceptibility



$$\delta Q = C_y \cdot \delta T \quad (5.54)$$

Def. Heat Capacity

$$C_y \equiv \left(\frac{\delta Q}{\delta T}\right)_y \quad (5.55)$$

$$\longrightarrow C_y = C_y(T, y) : \text{ext. parameter} \quad (5.56)$$

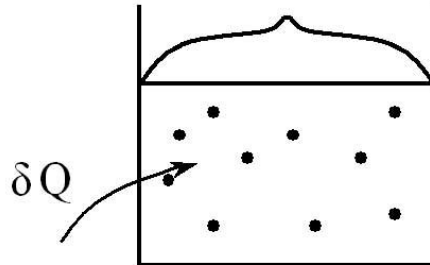
Ext. \longrightarrow Int.

$$\frac{C_y}{N}, \frac{C_y}{M}, \frac{C_y}{\nu} : \text{specific heat per ptl., mass, mole} \quad (5.57)$$

e. g. 1 Gas system

$$y = V, P \quad (5.58)$$

(1) C_V



$$dU = \delta Q - \delta W = T dS - P dV \quad (5.59)$$

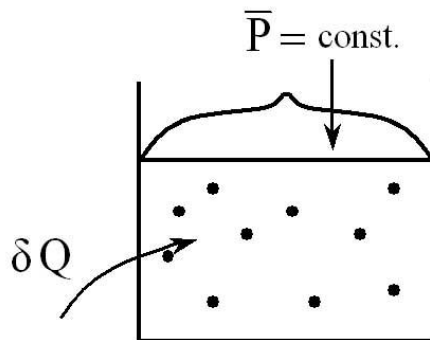
$$\downarrow V = \text{const}$$

$$dU = \delta Q \quad (5.60)$$

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V = T \cdot \left(\frac{\partial S}{\partial T} \right)_V \quad (5.61)$$

$$= -T \cdot \left(\frac{\partial^2 F}{\partial T^2} \right)_V \quad (5.62)$$

(2) C_P



$$\delta Q = dU + P \cdot dV = dH \quad (5.63)$$

$$dH = T \cdot dS + V \cdot dP \quad (5.64)$$

One can expect that $C_P > C_V$

$$C_P = \left(\frac{\delta Q}{\delta T} \right)_P = \left(\frac{\partial H}{\partial T} \right)_P = T \cdot \left(\frac{\partial S}{\partial T} \right)_P \quad (5.65)$$

$$= T \cdot \left[\left(\frac{\partial S}{\partial T} \right)_V + \left(\frac{\partial S}{\partial V} \right)_T \cdot \left(\frac{\partial V}{\partial T} \right)_P \right] \quad (5.66)$$

$$= C_V + T \cdot \left(\frac{\partial P}{\partial T} \right)_V \cdot \left(\frac{\partial V}{\partial T} \right)_P \quad (5.67)$$

$$\therefore C_P - C_V = T \cdot \left(\frac{\partial P}{\partial T} \right)_V \cdot \left(\frac{\partial V}{\partial T} \right)_P \quad (5.68)$$

$$V = V(T, P) \quad (5.69)$$

$$dV = \left(\frac{\partial V}{\partial T} \right)_P \cdot dT + \left(\frac{\partial V}{\partial P} \right)_T \cdot dP = 0 \quad (5.70)$$

$$\therefore \left(\frac{\partial P}{\partial T} \right)_V = - \frac{\left(\frac{\partial V}{\partial T} \right)_P}{\left(\frac{\partial V}{\partial P} \right)_T} : \text{Euler's relation} \quad (5.71)$$

$$\therefore C_P - C_V = -T \cdot \frac{\left[\left(\frac{\partial V}{\partial T} \right)_P \right]^2}{\left(\frac{\partial V}{\partial P} \right)_T} \quad (5.72)$$

Def. Coefficient of Volume Expansion

$$\alpha \equiv \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_P \quad (5.73)$$

Def. Isothermal Compressibility

$$\kappa = - \frac{1}{V} \left(\frac{\partial V}{\partial P} \right)_T \quad (5.74)$$

$$\therefore C_P - C_V = T \cdot V \cdot \frac{\alpha^2}{\kappa} > 0 \quad (5.75)$$

(3) Limiting Behavior of C_y ($T \rightarrow 0$)

$$T \longrightarrow 0^+, E = E_0 \quad (5.76)$$

$$S = k_B \ln \Omega(E_0) \quad (5.77)$$

$$= S_0 : \text{indep. of all parameters} \quad (5.78)$$

$$C_y = T \cdot \left(\frac{\partial S}{\partial T} \right)_y \rightarrow \left(\frac{\partial S}{\partial T} \right)_y = \frac{C_y}{T}; y = V \& P \quad (5.79)$$

$$\therefore S(T) - S(0) = \int_0^T \frac{C_y(T')}{T'} \cdot dT' \quad (5.80)$$

$$\therefore \text{As } T \rightarrow 0, C_y(T) \rightarrow 0. \quad (5.81)$$

$$(\because \text{ to guarantee the proper convergence of the integral}) \quad (5.82)$$

e. g. 2 Magnetic system

$$dU = T dS - M dB \quad (5.83)$$

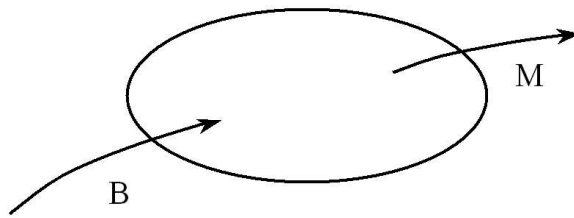
$$dF = -S dT - M dB, S = - \left(\frac{\partial F}{\partial T} \right)_B \ \& \ M = - \left(\frac{\partial F}{\partial B} \right)_T \quad (5.84)$$

• Specific heat

$$C_B = \left(\frac{\partial U}{\partial T} \right)_B = T \cdot \left(\frac{\partial S}{\partial T} \right)_B = - \left(\frac{\partial^2 F}{\partial T^2} \right)_B \quad (5.85)$$

• Magnetic Susceptibility

$$\chi = \left(\frac{\partial M}{\partial B} \right)_T = \left(\frac{\partial^2 F}{\partial B^2} \right)_T \quad (5.86)$$



e. g. 1 Ideal Gas

$$S = N k_B \left[\ln V + \frac{3}{2} \ln U \right] + S_0 \quad (5.87)$$

↓

$$U(S, V) = e^{\frac{2(S-S_0)}{3Nk_B}} \cdot V^{-\frac{2}{3}} \quad (5.88)$$

$$\frac{\partial U}{\partial S} = \frac{2}{3N k_B} \cdot U = T \longrightarrow U = \frac{3}{2} N k_B T \quad (5.89)$$

$$\frac{\partial U}{\partial V} = -\frac{2}{3} U = -P \longrightarrow P V = N k_B T \quad (5.90)$$

$$C_V = \left(\frac{\partial U}{\partial T} \right)_V = \frac{3}{2} N k_B \longrightarrow \frac{C_V}{N} = \frac{3}{2} k_B \quad (5.91)$$

$$H(S, P) = U(S, V) + P \cdot V \quad (5.92)$$

↓

$$H = \frac{5}{2} N k_B T \quad (5.93)$$

$$\therefore C_P = \left(\frac{\partial H}{\partial T} \right)_P = \frac{5}{2} N k_B \longrightarrow \frac{C_P}{N} = \frac{5}{2} k_B \quad (5.94)$$

$$\therefore C_P - C_V = N k_B \quad (5.95)$$

- Coefficient of volume expansion:

$$\alpha = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_P = \frac{1}{T} \quad (5.96)$$

- Isothermal compressibility

$$\kappa = -\frac{1}{V} \left(\frac{\partial V}{\partial P} \right)_T = \frac{1}{P} \quad (5.97)$$

e. g. 2 Magnetic System

$$S(U, B) = -N k_B [p_+ \ln p_+ + p_- \ln p_-] \quad (5.98)$$

$$p_+ = \frac{1}{2} \left[1 + \frac{U}{N\varepsilon} \right], \quad p_- = \frac{1}{2} \left[1 - \frac{U}{N\varepsilon} \right] \quad (5.99)$$

$$\varepsilon = -\mu B \quad (5.100)$$

$$\frac{\partial S}{\partial U} = \frac{1}{T} \longrightarrow U(T, B) = -N \cdot \varepsilon \cdot \tanh \frac{\varepsilon}{k_B T} \quad (5.101)$$

$$C_B = \left(\frac{\partial U}{\partial T} \right)_B = \frac{N\varepsilon^2}{k_B T^2} \cdot \operatorname{sech}^2 \frac{\varepsilon}{k_B T} \quad (5.102)$$

$$\frac{\partial S}{\partial B} = \frac{M}{T} \longrightarrow M = N\mu \cdot \tanh \left(\frac{\mu B}{k_B T} \right) \quad (5.103)$$

$$\chi_T = \left(\frac{\partial M}{\partial B} \right)_T = \frac{N\mu^2}{k_B T} \operatorname{sech}^2 \left(\frac{\mu B}{k_B T} \right) \propto \frac{1}{T} \quad (5.104)$$

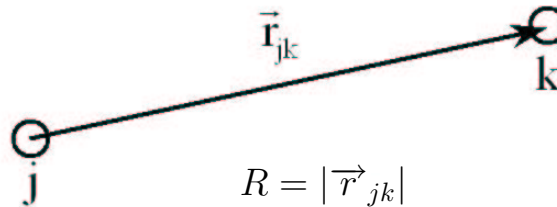
5.6 Van der Waals Gas

Non-ideal gas \longrightarrow interacting gas

$$H = T + U, \quad T = \sum_{i=1}^N \frac{\vec{p}_i^2}{2m} \quad (5.105)$$

$U?$

“pair interaction” u_{ik}



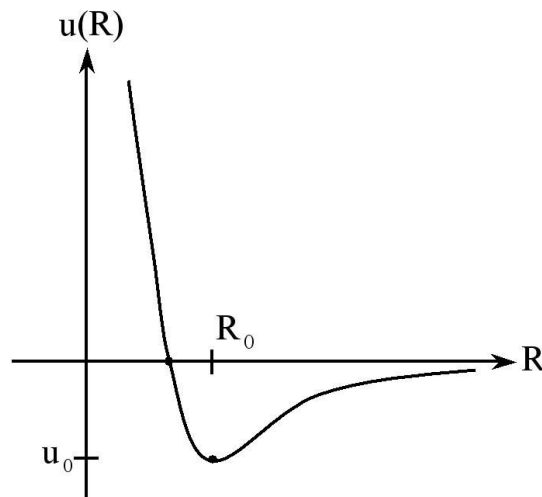
$$R = |\vec{r}_{jk}| \quad (5.106)$$

“Lennard-Jones potential” (semi-empirical pot.)

$$u(R) = u_0 \left[\left(\frac{R_0}{R} \right)^{12} - 2 \left(\frac{R_0}{R} \right)^6 \right] \quad (5.107)$$

u_0 & R_0 (empirical parameters):

constants depending on the molecules or atoms under consideration



long-range attractive part

+

short-range repulsive part

$$U = \frac{1}{2} \sum_{j, k=1(j \neq k)}^N u_{jk} \quad (5.108)$$

- (Empirical) Eq. of State

$$\left(p + \underbrace{a \cdot \frac{N^2}{V^2}}_{\text{long-range attraction}} \right) \cdot \left(V - \underbrace{b \cdot N}_{\text{short-range repulsion}} \right) = N \cdot k_B T \quad (5.109)$$

a, b : empirical positive parameters depending on the molecules or atoms

- Virial expansion

$$p + an^2 = nk_B T (1 - bn)^{-1} \quad (5.110)$$

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$$p = nk_B T [1 + bn + b^2 n^2 + \dots] - an^2 \quad (5.111)$$

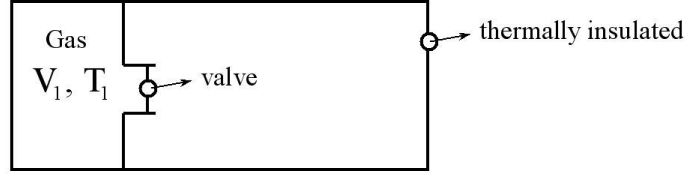
$$= nk_B T \left[1 + \left(b - \frac{a}{k_B T} \right) n + \dots \right] \quad (5.112)$$

Generally, $p = k_B T [n + B_2(T) \cdot n^2 + B_3(T) \cdot n^3 + \dots]$

$$B_2, B_3, \dots : \text{virial coefficients} \quad (5.113)$$

$$\text{Van der Waals} \longrightarrow B_2(T) = b - \frac{a}{k_B T} \quad (5.114)$$

5.7 Free Expansion



$$\text{Initial state: } V_1, T_1 \quad (5.115)$$

↓ valve: open

$$\text{Final state: } V_2, T_2(?) \quad (5.116)$$

$$Q = 0, W = 0 \longrightarrow U = \text{const.} \quad (5.117)$$

$$\left(\frac{\partial T}{\partial V} \right)_U ? \quad (5.118)$$

$$U = U(T, V) \rightarrow dU = \left(\frac{\partial U}{\partial V} \right)_T \cdot dV + \left(\frac{\partial U}{\partial T} \right)_V \cdot dT \quad (5.119)$$

$$= 0 \quad (5.120)$$

$$\therefore \left(\frac{\partial T}{\partial V} \right)_U = - \frac{\left(\frac{\partial U}{\partial V} \right)_T}{\left(\frac{\partial U}{\partial T} \right)_V} = - \frac{1}{C_V} \cdot \left(\frac{\partial U}{\partial V} \right)_T \quad (5.121)$$

$$U = U(S, V) \quad (5.122)$$

$$\left(\frac{\partial U}{\partial V} \right)_T = \left(\frac{\partial U}{\partial V} \right)_S + \left(\frac{\partial U}{\partial S} \right)_V \cdot \left(\frac{\partial S}{\partial V} \right)_T \quad (5.123)$$

$$= -p + T \cdot \left(\frac{\partial p}{\partial T} \right)_V \quad (5.124)$$

$$\left(\frac{\partial p}{\partial T} \right)_V = - \frac{\left(\frac{\partial V}{\partial T} \right)_p}{\left(\frac{\partial V}{\partial p} \right)_T} = \frac{\alpha_p}{\kappa_T} \quad (5.125)$$

$$\therefore \left(\frac{\partial U}{\partial V} \right)_T = -p + \frac{T \cdot \alpha_p}{\kappa_T} \quad (5.126)$$

$$\therefore \left(\frac{\partial T}{\partial V} \right)_U = \frac{1}{C_V} \left[p - T \cdot \frac{\alpha_p}{\kappa_T} \right] \quad (5.127)$$

e. g. 1 Ideal Gas

$$\left(\frac{\partial T}{\partial V}\right)_U = 0 \quad T_2 = T_1 \quad (5.128)$$

e. g. 2 Van der Waals Gas

$$p \simeq k_B T [n + B_2(T)n^2] = nk_B T (1 + B_2 \cdot n) \quad (5.129)$$

$$B_2(T) = b - \frac{a}{k_B T} \quad (5.130)$$

$$V \simeq \frac{Nk_B T}{p} (1 + B_2 \cdot n) \quad (5.131)$$

$$= N \left(\frac{k_B T}{p} + B_2 \cdot \frac{nk_B T}{p} \right) \quad (5.132)$$

$$\simeq N \cdot \left(\frac{k_B T}{p} + B_2 \right) \quad (5.133)$$

$$\alpha_p = \frac{1}{V} \left(\frac{\partial V}{\partial T} \right)_p \simeq \frac{N}{V} \left(\frac{k_B}{p} + \frac{\partial B_2}{\partial T} \right) \quad (5.134)$$

$$\kappa_T = -\frac{1}{V} \left(\frac{\partial V}{\partial p} \right)_T \simeq \frac{N}{V} \cdot \frac{k_B T}{p^2} \quad (5.135)$$

$$\therefore \left(\frac{\partial T}{\partial V} \right)_U = \frac{1}{C_V} \left(-\frac{p^2}{k_B} \right) \frac{\partial B_2}{\partial T} \quad (5.136)$$

$$\frac{\partial B_2}{\partial T} = \frac{a}{k_B T^2} > 0 \quad (5.137)$$

$$\therefore \left(\frac{\partial T}{\partial V} \right)_U < 0 \quad (5.138)$$

$$\left(\frac{\partial T}{\partial V} \right)_U = -\frac{1}{C_V} \cdot a \cdot \left(\frac{p}{k_B T} \right)^2 \quad (5.139)$$

$$\simeq -\frac{a}{C_V} \cdot n^2 < 0 \quad (5.140)$$

$$T_2 < T_1 \quad (5.141)$$

5.8 Adiabatic Expansion

Quasi-static adiabatic process: $\Delta S = 0$

$$\left(\frac{\partial T}{\partial V}\right)_S \quad ? \quad (5.142)$$

$$S = S(T, V) \quad (5.143)$$

$$dS = \left(\frac{\partial S}{\partial T}\right)_V \cdot dT + \left(\frac{\partial S}{\partial V}\right)_T \cdot dV = 0 \quad (5.144)$$

$$\therefore \left(\frac{\partial T}{\partial V}\right)_S = -\frac{\left(\frac{\partial S}{\partial V}\right)_T}{\left(\frac{\partial S}{\partial T}\right)_V} \quad (5.145)$$

$$= -\frac{T}{C_V} \cdot \left(\frac{\partial P}{\partial T}\right)_V \quad (5.146)$$

$$= -\frac{T}{C_V} \cdot \frac{\alpha_P}{\kappa_T} \quad (5.147)$$

e. g. Ideal Gas

$$\left(\frac{\partial T}{\partial V}\right)_S = -\frac{T}{C_V} \cdot \left(\frac{\partial P}{\partial T}\right)_V, \quad P \cdot V = Nk_B T, \quad C_V = \frac{3}{2}Nk_B \quad (5.148)$$

$$= -\frac{T}{C_V} \cdot \frac{Nk_B}{V} \quad (5.149)$$

$$= -\frac{3}{2} \cdot \frac{T}{V} \quad (5.150)$$

$$\therefore \frac{1}{T} dT = -\frac{2}{3} \cdot \frac{1}{V} dV \quad (5.151)$$

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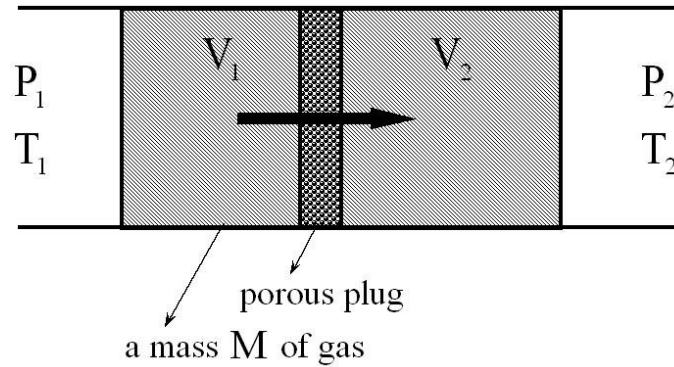
$$T_f \cdot V_f^{\frac{2}{3}} = T_i \cdot V_i^{\frac{2}{3}} \quad (5.152)$$

$$\gamma \equiv \frac{C_P}{C_V} = \frac{5}{3} \longrightarrow T \cdot V^{\gamma-1} = \text{const.} \quad (5.153)$$

↓

$$P \cdot V^\gamma = \text{const.} \quad (5.154)$$

5.9 Throttling Process (Joule-Thomson process)



thermally insulated P_1 & P_2 : const. $P_1 > P_2$

The gas to the left of the mass $M \rightarrow$ mass $M = P_1 \cdot V_1$

Mass $M \rightarrow$ the gas to the right of it = $P_2 \cdot V_2$

$$Q = 0, \quad (5.155)$$

$$W = P_2 \cdot V_2 - P_1 \cdot V_1 \quad (5.156)$$

$$\therefore \Delta U = -W = P_1 \cdot V_1 - P_2 \cdot V_2 \quad (5.157)$$

$$\therefore U_1 + P_1 \cdot V_1 = U_2 + P_2 V_2 \quad (5.158)$$

$$\therefore H_1 = H_2 \longrightarrow H = \text{const} \quad (5.159)$$

$$H(T_1, P_1) = H(T_2, P_2) \quad (5.160)$$

$$\text{J. T. coefficient } \mu \equiv \left(\frac{\partial T}{\partial P} \right)_H \quad (5.161)$$

$$H = H(T, P) \longrightarrow dH = \left(\frac{\partial H}{\partial T} \right)_P \cdot dT + \left(\frac{\partial H}{\partial P} \right)_T \cdot dP = 0 \quad (5.162)$$

$$\therefore \mu = \left(\frac{\partial T}{\partial P} \right)_H = - \frac{\left(\frac{\partial H}{\partial P} \right)_T}{\left(\frac{\partial H}{\partial T} \right)_P} = - \frac{1}{C_P} \cdot \left(\frac{\partial H}{\partial P} \right)_T \quad (5.163)$$

$$H = H(S, P) \longrightarrow dH = T dS + V dP \quad (5.164)$$

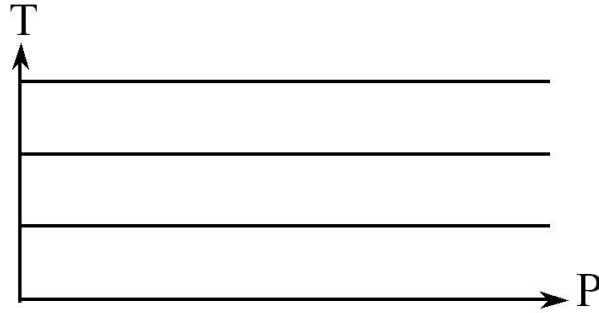
$$\left(\frac{\partial H}{\partial P}\right)_T = \left(\frac{\partial H}{\partial P}\right)_S + \left(\frac{\partial H}{\partial S}\right)_P \cdot \left(\frac{\partial S}{\partial P}\right)_T \quad (5.165)$$

$$= V - T \cdot \left(\frac{\partial V}{\partial T}\right)_P \quad (5.166)$$

$$\therefore \mu = \left(\frac{\partial T}{\partial P}\right)_H = \frac{V}{C_P}(T \cdot \alpha_P - 1) \quad (5.167)$$

e. g. 1 Ideal Gas

$$\alpha_P = \frac{1}{T} \longrightarrow \mu = 0 \text{ (no temp. change)} \quad (5.168)$$



$$H = \frac{5}{2}k_B T \quad \left(\frac{\partial T}{\partial P}\right)_H = 0 \quad (5.169)$$

e. g. 2 Van der Waals Gas

$$\alpha_P = \frac{1}{V} \left(\frac{\partial V}{\partial T}\right)_P = \frac{Nk_B}{P \cdot V} + \frac{N}{V} \frac{\partial B_2}{\partial T} \quad (5.170)$$

$$T \cdot \alpha_P - 1 = \frac{Nk_B T}{P \cdot V} + \frac{NT}{V} \frac{\partial B_2}{\partial T} - 1 \quad (5.171)$$

$$\frac{Nk_B T}{P \cdot V} = (1 + B_2 \cdot n)^{-1} \simeq 1 - B_2 \cdot n \quad (5.172)$$

$$\therefore T \cdot \alpha_P - 1 \simeq (T \cdot \frac{\partial B_2}{\partial T} - B_2) \cdot n \quad (5.173)$$

$$\frac{\partial B_2}{\partial T} = \frac{a}{k_B T^2} \quad (5.174)$$

$$\therefore \mu \simeq \frac{N}{C_P} (T \cdot \frac{\partial B_2}{\partial T} - B_2) = \frac{N}{C_P} \left(\frac{2a}{k_B T} - b \right) \quad (5.175)$$

For sufficiently low temp., $\frac{a}{k_B T}$ (attractive int.)

$> b$ (repulsive int.)

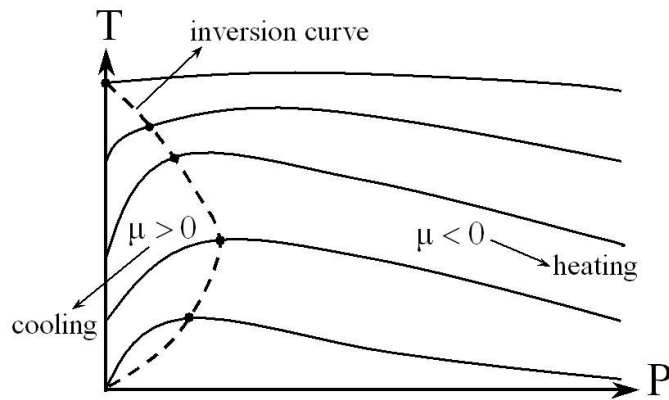
$$\therefore \mu > 0 \text{ (cooling)} \quad (5.176)$$

For sufficiently high temp., $\frac{a}{k_B T} < b$

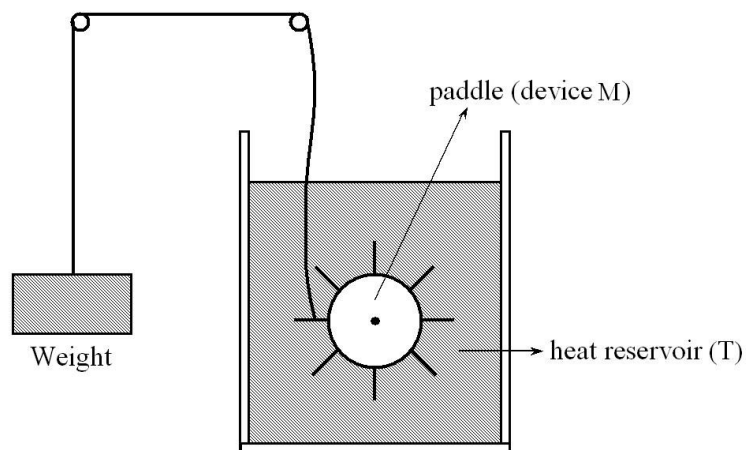
$$\therefore \mu < 0 \text{ (heating)} \quad (5.177)$$

For some intermediate temp., $\frac{a}{k_B T} = b$

$$\mu = 0 \quad (5.178)$$



5.10 Heat Engine

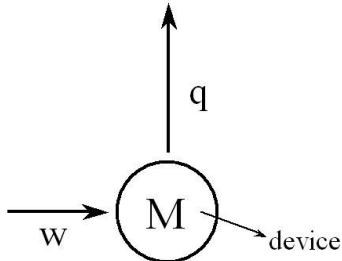
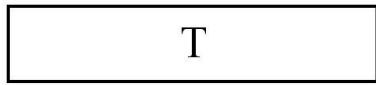


- Falling a weight

Conversion of mechanical work into heat



Transformation of order into randomness



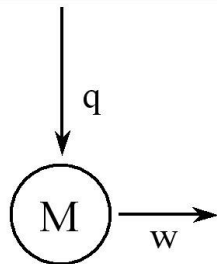
many degrees of freedom (random motion)

$$\Delta S = \frac{q}{T} > 0$$

(a single degree of freedom): ordered motion

- Lifting a weight

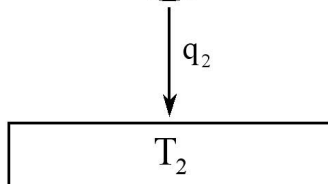
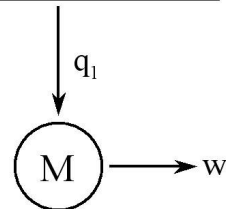
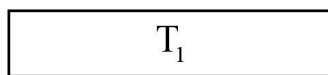
Conversion of heat into work (reverse process)



$$\Delta S = -\frac{q}{T} < 0$$

Perfect Engine for one cycle

Real Engine



$$T_1 > T_2$$

$$q_1 = W + q_2, \Delta S = -\frac{q_1}{T_1} + \frac{q_2}{T_2} \geq 0 \quad (5.179)$$

$$-\frac{q_1}{T_1} + \frac{q_1 - W}{T_2} \geq 0 \longrightarrow \frac{W}{T_2} \leq q_1 \left(\frac{1}{T_2} - \frac{1}{T_1} \right) \quad (5.180)$$

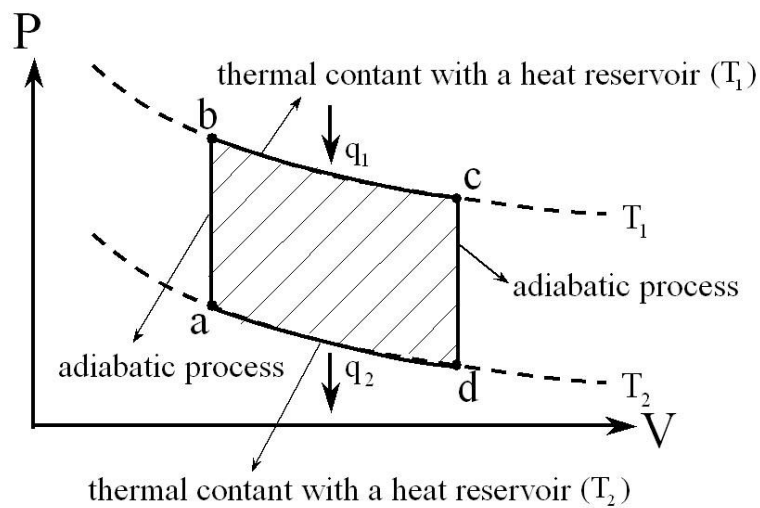
$$\therefore T_2 < T_1 \quad (5.181)$$

$$\text{Efficiency } \eta \equiv \frac{W}{q_1} \leq 1 - \frac{T_1}{T_2} = \frac{T_1 - T_2}{T_1} \quad (5.182)$$

= : quasi-static process (max. work)

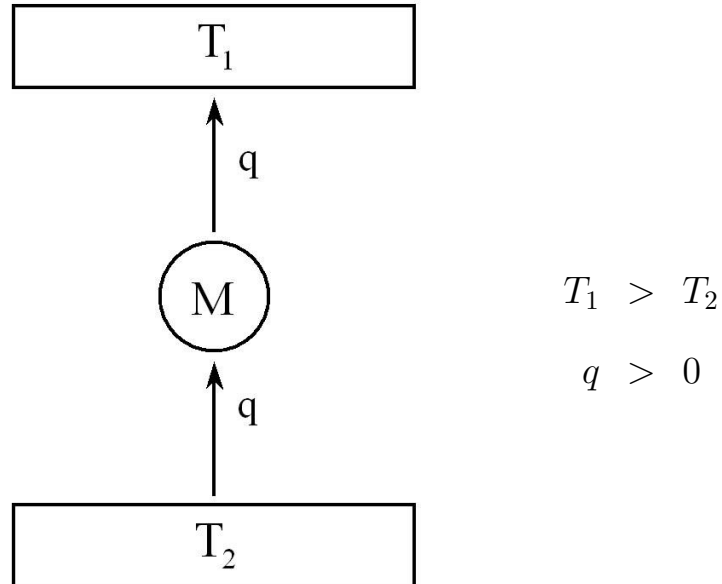
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e. g. Carnot Engine



5.11 Refrigerators

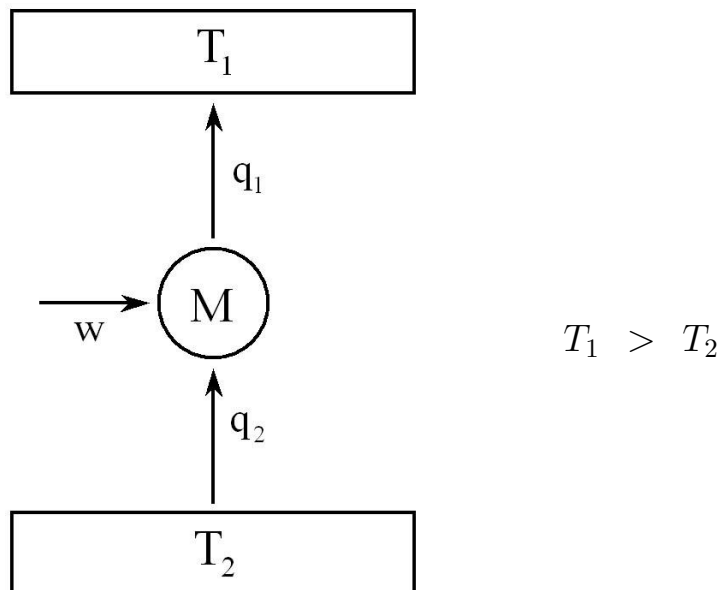
- a perfect refrigerator



$$\Delta S = q \left(\frac{1}{T_1} - \frac{1}{T_2} \right) < 0 \quad (5.183)$$

\therefore impossible

- a real refrigerator



$$\Delta S = \frac{q_1}{T_1} - \frac{q_2}{T_2} \geq 0 \longrightarrow \frac{q_2}{q_1} \leq \frac{T_2}{T_1} \quad (5.184)$$

$$\frac{q_2}{q_2 + W} \leq \frac{T_2}{T_1} \quad (5.185)$$

$$\downarrow$$
$$q_2 \leq W \left(\frac{T_2}{T_1 - T_2} \right) \quad (5.186)$$

Coefficient of performance of a refrigerate

$$K \equiv \frac{q_2}{W} \leq \frac{T_2}{T_1 - T_2} \quad (5.187)$$