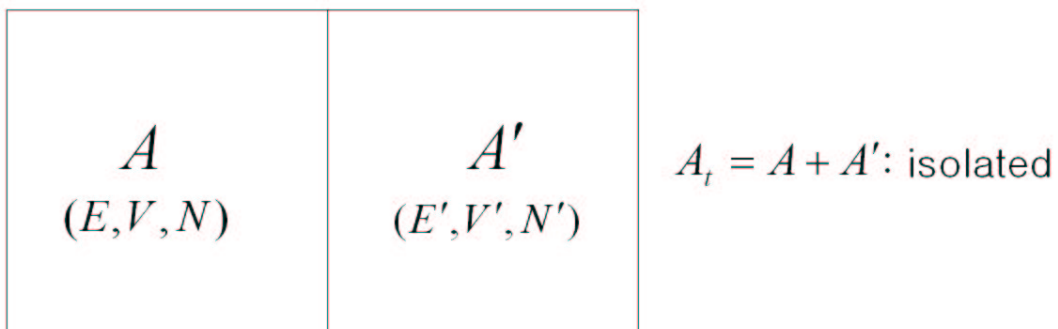


# Chapter 4

## Grandcanonical Ensemble

### 4.1 Chemical Diffusive Equilibrium and Chemical Potential



$$E + E' = E_t : \text{fixed.} \tag{4.1}$$

$$V + V' = V_t : \text{fixed.} \tag{4.2}$$

$$N + N' = N_t : \text{fixed.} \tag{4.3}$$

$$S_t(E, V, N) = S(E, V, N) + S'(E', V', N') \tag{4.4}$$

Equilibrium Conditions:  $S_t$ : Max

$$\frac{\partial S_t}{\partial E} = 0 \longrightarrow \frac{\partial S}{\partial E} = \frac{\partial S'}{\partial E'} \tag{4.5}$$

$$\longrightarrow T = T' : \text{Thermal Equilibrium} \tag{4.6}$$

$$\frac{\partial S_t}{\partial V} = 0 \longrightarrow \frac{\partial S}{\partial V} = \frac{\partial S'}{\partial V'} \tag{4.7}$$

$$\longrightarrow P = P' : \text{Mechanical Equilibrium} \tag{4.8}$$

$$\frac{\partial S_t}{\partial N} = 0 \longrightarrow \frac{\partial S}{\partial N} = \frac{\partial S'}{\partial N'} \tag{4.9}$$

• Def. Chemical Potential

$$\mu \equiv -T \cdot \frac{\partial S}{\partial N} \tag{4.10}$$

Note:  $\frac{\partial S}{\partial N}$  값이 큰 계가 ( $\mu$  값이 작은 계) 입자를 획득해야, 다른 계가 잃은 것보다 더 많은 엔트로피를 얻게 된다.: 입자흐름 ( $\mu$  값이 큰 계  $\rightarrow$   $\mu$  값이 작은 계 )

$$\mu = \mu' : \text{Chemical Diffusive Equilibrium} \tag{4.11}$$

$$S = S(E, V, N) \tag{4.12}$$

$$\rightarrow dS = \frac{1}{T}dE + \frac{P}{T}dV - \frac{\mu}{T}dN \tag{4.13}$$

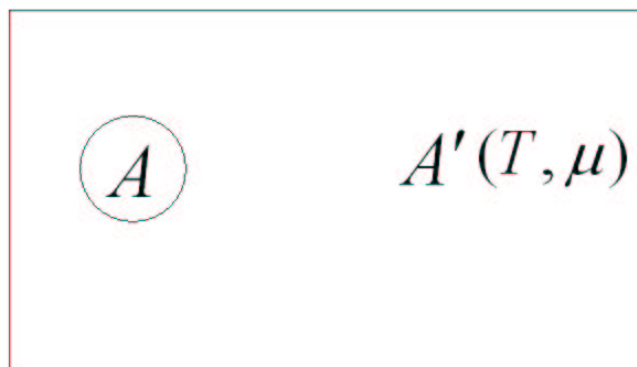
$$\rightarrow dE = TdS - PdV + \mu dN : \text{Thermodynamic 1st Law} \tag{4.14}$$

$$[-\mu dN : \text{chemical work}] \tag{4.15}$$

$$E = E(S, V, N), H = H(S, P, N), F = F(T, V, N), \tag{4.16}$$

$$G = G(T, P, N) : \text{Thermodynamic Potentials} \tag{4.17}$$

4.2 Grand Canonical Ensemble



$A_t = A$  (system) +  $A'$  (ptl-energy reservoir) : isolated

$N + N' = N_t = \text{fixed}, E + E' = E_t = \text{fixed}$

$$\frac{N}{N_t} \ll 1, \quad \frac{E}{E_t} \ll 1$$

• **Equilibrium Conditions**

$$P(E, N) = C \Omega_t(E, N) = C \Omega(E, N) \cdot \Omega'(E_t - E, N_t - N) \quad (4.18)$$

What is the most probable state?

$$P(E, N) : \max \iff \Omega_t : \max \quad (4.19)$$

$$S_t = k_B \ln \Omega_t : \max \quad (4.20)$$

$$= S(E, N) + S'(E_t - E, N_t - N) \quad (4.21)$$

$$\frac{\partial S_t}{\partial E} = \frac{\partial S}{\partial E} - \frac{\partial S'}{\partial E'} = 0 \longrightarrow \frac{\partial S}{\partial E} = \frac{\partial S'}{\partial E'} \quad (4.22)$$

$$\frac{\partial S_t}{\partial N} = \frac{\partial S}{\partial N} - \frac{\partial S'}{\partial N'} = 0 \longrightarrow \frac{\partial S}{\partial N} = \frac{\partial S'}{\partial N'} \quad (4.23)$$

$$\frac{\partial S'}{\partial E'} = \frac{1}{T'}, \quad \frac{\partial S'}{\partial N'} = -\frac{\mu'}{T'} \quad (4.24)$$

$$\therefore \text{평행조건} : T = T', \quad \mu = \mu' \quad (4.25)$$

$$A' : \text{reservoir} \longrightarrow \frac{E}{E_t} \ll 1, \quad \frac{N}{N_t} \ll 1$$

$$\therefore T' = \text{const.}, \quad \mu' = \text{const.} \quad (4.26)$$

$\therefore$  우리가 관심을 두는 계 A의 온도와 화학 포텐셜이  $T'$ 와  $\mu'$ 이 되면 평행 상태에 도달함.

• **Probability Distribution in Equilibrium**

가정: 계 A의 입자가  $N$ 개이고  $r$ 상태에 있다고 하자.

$$\Omega_t(N, E_r(N)) = 1 \cdot \Omega'(N_t - N, E_t - E_r(N)) \quad (4.27)$$

$$P_r(N) = C \cdot \Omega'(E_t - E_r(N), N_t - N) \quad (4.28)$$

$$\therefore \ln P_r(N) = \ln C + \ln \Omega'(E_t - E_r, N_t - N) \quad (4.29)$$

$$\approx \text{const.} + \left. \frac{\partial \ln \Omega'}{\partial E'} \right|_{(E_t, N_t)} \cdot (-E_r) + \left. \frac{\partial \ln \Omega'}{\partial N'} \right|_{(E_t, N_t)} \cdot (-N) \quad (4.30)$$

$$\therefore \beta \equiv \left. \frac{\partial \ln \Omega'}{\partial E'} \right|_{(E_t, N_t)} = \frac{1}{k_B} \frac{\partial S'}{\partial E'} = \frac{1}{k_B T} \quad (4.31)$$

$$\alpha \equiv \left. \frac{\partial \ln \Omega'}{\partial N'} \right|_{(E_t, N_t)} = \frac{1}{k_B} \frac{\partial S'}{\partial N'} = -\frac{\mu}{k_B T} = -\beta \mu \quad (4.32)$$

$$\therefore P_r(N) = C \cdot e^{-\beta E_r - \alpha N} = C \cdot e^{-\beta(E_r - \mu N)} \quad [e^{-\beta(E_r - \mu N)} : \text{Gibbs factor}] \quad (4.33)$$

$$\longrightarrow \sum_{N=0}^{\infty} \sum_r P_r(N) = C \cdot \sum_{N=0}^{\infty} e^{-\alpha N} \sum_r e^{-\beta E_r(N)} \quad (4.34)$$

$$= C \cdot \sum_{N=0}^{\infty} e^{-\alpha N} \cdot Z(N) = 1 \quad (4.35)$$

### • Def. Grand Partition Function

$$Q_G = \sum_{N=0}^{\infty} e^{\beta \mu N} Z(N); \quad Z(N) = \sum_r e^{-\beta E_r(N)} \quad (4.36)$$

$$\rightarrow C = 1/Q_G, \quad P_r(N) = e^{-\beta(E_r - \mu N)}/Q_G \quad (4.37)$$

### • Grand Potential

$$Q_G = \sum_{N=0}^{\infty} e^{\beta \mu N} \cdot Z(N); \quad Z(N) = \sum_r e^{-\beta E_r(N)} \quad (4.38)$$

$z \equiv e^{\beta \mu}$ : absolute activity or fugacity of the system

$$Q_G = \sum_{N=0}^{\infty} z^N \cdot Z(T, V, N) \quad (4.39)$$

$$\Omega(\text{grand potential}) = \Omega(T, V, \mu) \quad (4.40)$$

$$\equiv F(T, V, \bar{N}) - \mu \bar{N} \quad (4.41)$$

$$d\Omega = dF - d(\mu\bar{N}) = -SdT - PdV - \bar{N}d\mu \quad (4.42)$$

$$S = -\left(\frac{\partial\Omega}{\partial T}\right)_{V,\mu}, \quad P = -\left(\frac{\partial\Omega}{\partial V}\right)_{T,\mu}, \quad \bar{N} = -\left(\frac{\partial\Omega}{\partial\mu}\right)_{T,V} \quad (4.43)$$

$$S = -k_B \sum_{(N,r)} P_r(N) \cdot \ln P_r(N) \quad (4.44)$$

$$\downarrow P_r(N) = \frac{e^{-\beta(E_r(N)) - \mu N}}{Q_G} \quad (4.45)$$

$$S = -k_B \sum_{(N,r)} \frac{e^{-\beta[E_r(N) - \mu N]}}{Q_G} \cdot [-\beta E_r(N) + \beta\mu N - \ln Q_G] \quad (4.46)$$

$$= \beta k_B \bar{E} - \beta\mu k_B \bar{N} + k_B \ln Q_G \quad (4.47)$$

$$S = \beta k_B \bar{E} - \beta\mu k_B \bar{N} + k_B \ln Q_G \quad (4.48)$$

↓

$$F = \bar{E} - TS = \mu\bar{N} - k_B T \ln Q_G \quad (4.49)$$

↓

$$\Omega = F - \mu\bar{N} = -k_B T \ln Q_G \quad (4.50)$$

★ 열역학적 극한에서는, 에너지와 입자수 요동 무시 가능

$$\longrightarrow Q_G \simeq e^{\beta\mu\bar{N}} \cdot Z(\bar{N}) \longrightarrow -k_B T \ln Q_G \quad (4.51)$$

$$= -k_B T \ln Z(\bar{N}) - \mu\bar{N} \quad (4.52)$$

$$= F - \mu\bar{N} \quad (4.53)$$

$$= \Omega(T, V, \mu) \quad (4.54)$$

## • Summary

1) Calculate  $Q_G$

$$2) \Omega(T, V, \mu) = -k_B T \ln Q_G(T, V, \mu) \quad (4.55)$$

$$F(T, V, \bar{N}) = \Omega + \bar{N}\mu; \quad (4.56)$$

$$\begin{aligned} &\downarrow \\ \bar{N} &= - \left( \frac{\partial \Omega}{\partial \mu} \right)_{T, V} \end{aligned} \quad (4.57)$$

$$\begin{aligned} &\downarrow \\ \bar{N} &= \bar{N}(T, V, \mu) \end{aligned} \quad (4.58)$$

$$\begin{aligned} &\downarrow \\ \mu &= \mu(T, V, \bar{N}) \end{aligned} \quad (4.59)$$

3) Other thermodynamic quantities

$$S = - \left( \frac{\partial \Omega}{\partial T} \right)_{V, \mu}, \quad - \left( \frac{\partial F}{\partial T} \right)_{V, \bar{N}} \quad (4.60)$$

$$P = - \left( \frac{\partial \Omega}{\partial V} \right)_{T, \mu}, \quad - \left( \frac{\partial F}{\partial V} \right)_{T, \bar{N}} \quad (4.61)$$

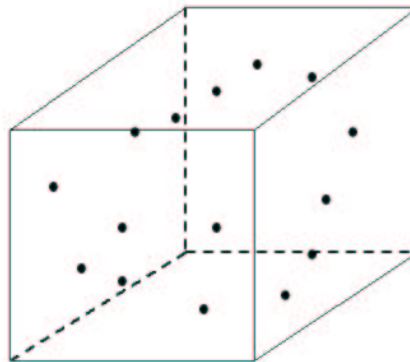
$$\bar{N} = - \left( \frac{\partial \Omega}{\partial \mu} \right)_{T, V} \quad (4.62)$$

$$\mu = \left( \frac{\partial F}{\partial \bar{N}} \right)_{T, V} \quad (4.63)$$

$$U = F - TS, \quad C_V = \left( \frac{\partial V}{\partial T} \right)_{V, \bar{N}} \quad (4.64)$$

### 4.3 Ideal Gas

e.g. Classical ideal gas



$$\mathcal{H} = \sum_{i=1}^{3N} \frac{p_i^2}{2m} \quad (4.65)$$

$$Q_G = \sum_{N=0}^{\infty} z^N \cdot Z_N, \quad z = e^{\beta\mu} \quad (4.66)$$

$$Z_N = \frac{1}{N!} \cdot \int \frac{\prod_{i=1}^{3N} dq_i \cdot dp_i}{h_0^{3N}} \cdot e^{-\beta \sum_{i=1}^{3N} \frac{p_i^2}{2m}} \quad (4.67)$$

$$= \frac{1}{N!} \cdot q^N, \quad q = \int \frac{\prod_{i=1}^3 dq_i \cdot dp_i}{h_0^3} \cdot e^{-\beta \sum_{i=1}^3 \frac{p_i^2}{2m}} \quad (4.68)$$

$$= V/\lambda_T^3, \quad \lambda_T = \sqrt{\frac{h_0^2}{2\pi m k_B T}} \quad (4.69)$$

$$\therefore Q_G = \sum_{N=0}^{\infty} z^N \cdot q^N / N! = e^{zq} \quad \text{for } zq < 1 \quad (4.70)$$

$$\Omega(T, V, \mu) = -k_B T \cdot \ln Q_G = -k_B T \cdot zq \quad (4.71)$$

$$= -k_B T \cdot e^{\beta\mu} \cdot V/\lambda_T^3 \quad (4.72)$$

$$\bar{N} = -\frac{\partial \Omega}{\partial \mu} = e^{\beta\mu} \cdot V/\lambda_T^3 \longrightarrow \mu = k_B T \cdot \ln(\lambda_T^3 \cdot \bar{n}), \quad \bar{n} = \bar{N}/V \quad (4.73)$$

$$= k_B T \cdot \ln(\lambda_T^3 / (V/\bar{N})) \quad (4.74)$$

$$\text{Classical case: } \frac{V}{\bar{N}} \gg \lambda_T^3 \longrightarrow \mu < 0 \quad (4.75)$$

$$P = -\frac{\partial \Omega}{\partial V} = k_B T \cdot e^{\beta\mu} / \lambda_T^3 = k_B T \cdot \frac{\bar{N}}{V} \longrightarrow PV = \bar{N} k_B T \quad (4.76)$$

$$U = -\frac{\partial \ln Q_G(T, V, z)}{\partial \beta} = \frac{3}{2} k_B T \cdot zq = \frac{3}{2} \bar{N} k_B T \quad (4.77)$$

$$u = U/\bar{N} = \frac{3}{2} k_B T \longrightarrow C_V = \frac{\partial u}{\partial T} = \frac{3}{2} k_B \quad (4.78)$$